





# Direct Simulation for Wind Instrument Synthesis

DAFX 08



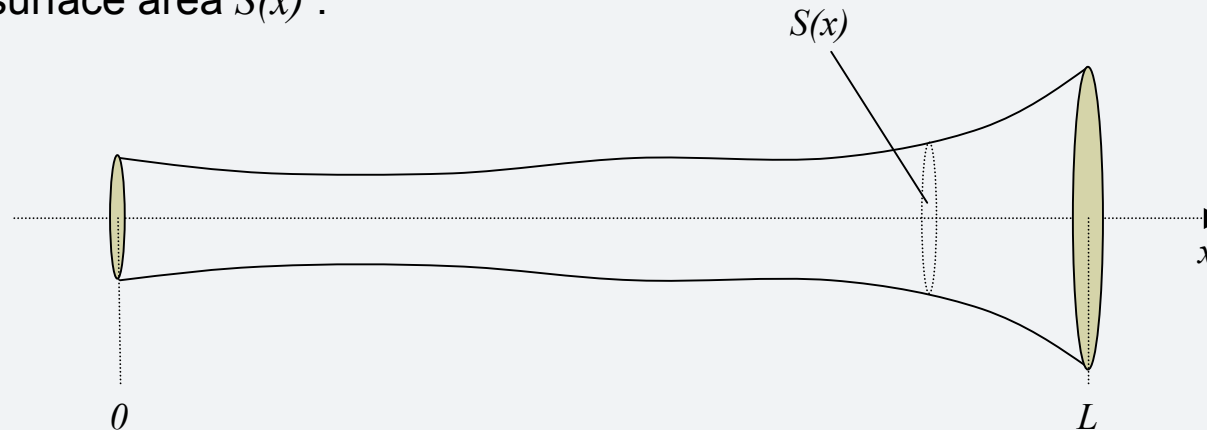
*Stefan Bilbao*

*Acoustics and Fluid Dynamics Group / Music  
University of Edinburgh*

-  **Webster's equation**
-  **Finite difference schemes**
-  **Efficiency, accuracy and stability**
-  **Sound examples: Single reed wind instruments**

# Webster's Equation

- Usual starting point for wind instrument models (and speech): an acoustic tube, surface area  $S(x)$  :



- Under various assumptions, velocity potential  $\Psi(x,t)$  satisfies:

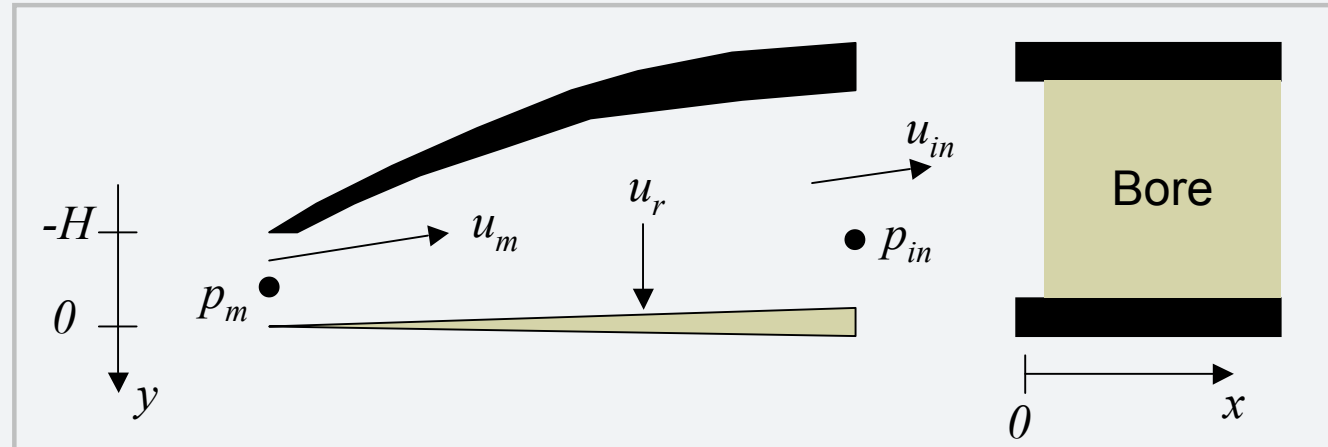
$$\Psi_{tt} = c^2 (S \Psi_x)_x$$

- $\Psi(x,t)$  related to pressure  $p(x,t)$  and volume velocity  $u(x,t)$  by:

$$p = \rho \Psi_t \qquad u = -S \Psi_x$$

# Single Reed Model

- A standard one-mass reed model:



- A driven oscillator:

$$\ddot{y} + g\dot{y} + \omega_0^2 y + \omega_1^{1+\alpha} \left( [-y - H]^+ \right)^\alpha = -a_1 p_\Delta$$

Linear oscillator terms      Collision term      Driving term

Mouthpiece pressure drop

$$p_\Delta = p_m - p_{in}$$

Flow nonlinearity

$$u_m = a_2 [y + H]^+ \sqrt{|p_\Delta|} \operatorname{sgn}(p_\Delta)$$

Flow conservation

$$u_{in} = u_m - u_r$$

Flow induced by reed

$$u_r = a_3 \dot{y}$$

Bore coupling

$$p_{in} = \rho \Psi_t(0, t)$$

$$u_{in} = -S(0) \Psi_x(0, t)$$

# Radiation Boundary Condition

- At the radiating end ( $x=L$ ), an approximate boundary condition is often given in impedance form:

$$P(s) = Z(s)U(s)$$

$$Z(s) = As - Bs^2$$

- Models inertial mass and loss.
- BUT: not positive real  $\rightarrow$  not passive.
- A better approximation (p.r., passive):

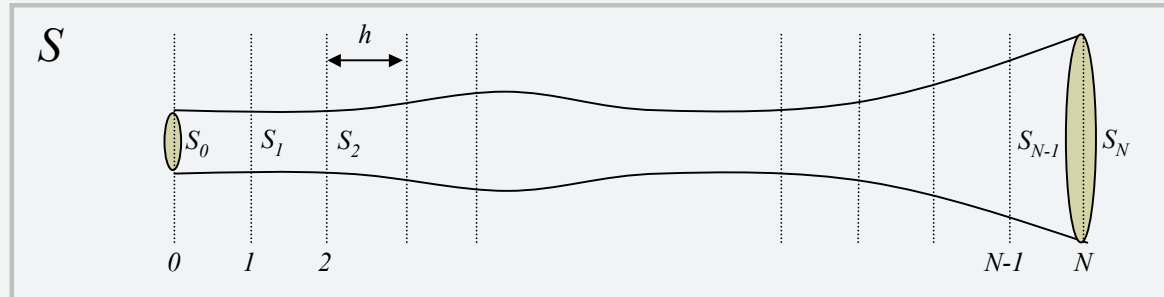
$$Z(s) = \frac{As}{1 + Bs / A}$$

- When converted to the time domain:

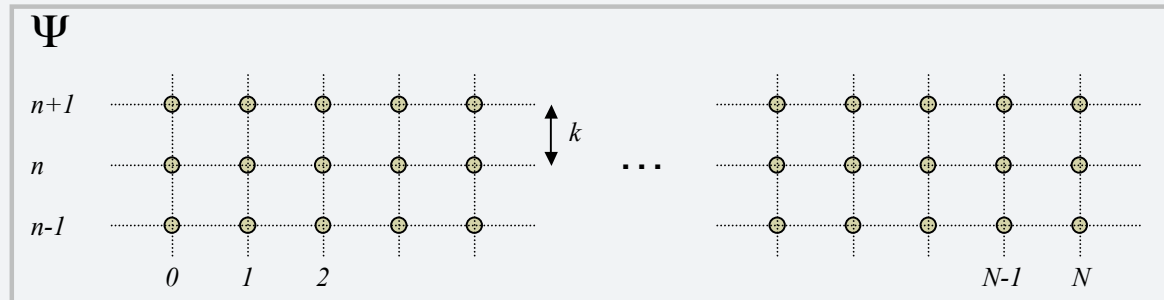
$$\Psi_x + q_1 \Psi_t + q_2 \Psi = 0 \quad \text{at} \quad x = L$$

# Finite Difference Scheme

- Sample bore profile at locations  
 $x = lh, l = 0, \dots, N$
- $h =$  grid spacing



- Introduce grid function  $\Psi$ , at locations  
 $x = lh, l = 0, \dots, N$   
 $t = nk, n = 0, \dots$
- $k =$  time step



- Here is one particular finite difference scheme (explicit, 2<sup>nd</sup> order accurate)

$$\Psi_l^{n+1} = 2\lambda^2 \frac{S_l + S_{l+1}}{S_{l+1} + 2S_l + S_{l-1}} \Psi_{l+1}^n + 2\lambda^2 \frac{S_l + S_{l-1}}{S_{l+1} + 2S_l + S_{l-1}} \Psi_{l-1}^n + 2(1 - \lambda^2) \Psi_l^n - \Psi_l^{n-1}$$

- Courant number  $\lambda$  defined as  $\lambda = ck/h$

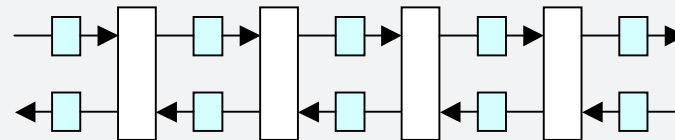
# Stability and Special Forms

- Can show (energy methods) that scheme is stable, over interior, when  
 $\lambda \leq 1$

- When  $\lambda = 1$ , scheme simplifies to:

$$\Psi_l^{n+1} = 2 \frac{S_l + S_{l+1}}{S_{l+1} + 2S_l + S_{l-1}} \Psi_{l+1}^n + 2 \frac{S_l + S_{l-1}}{S_{l+1} + 2S_l + S_{l-1}} \Psi_{l-1}^n - \Psi_l^{n-1}$$

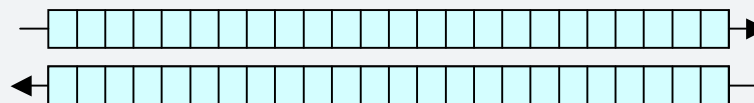
...equivalent to Kelly-Lochbaum scattering method



- When  $\lambda = 1$ , and  $S = \text{const.}$ , scheme simplifies further:

$$\Psi_l^{n+1} = \Psi_{l+1}^n + \Psi_{l-1}^n - \Psi_l^{n-1}$$

...equivalent to digital waveguide (exact integrator)



# Stability Condition and Tuning

- Stability condition requires  $\lambda \leq 1 \rightarrow h \geq ck$
- For simplicity, would like to choose an  $h$  which divides  $L$  evenly, i.e.,

$$L/h = N \quad \text{for integer } N$$

- Not possible for waveguide/Kelly-Lochbaum methods ---  $h=ck$ . Result: detuning, remedied using fractional delays.
- In an FD scheme, can choose  $h$  as one wishes. Result: very minor dispersion/loss of audio bandwidth. Numerical cutoff:

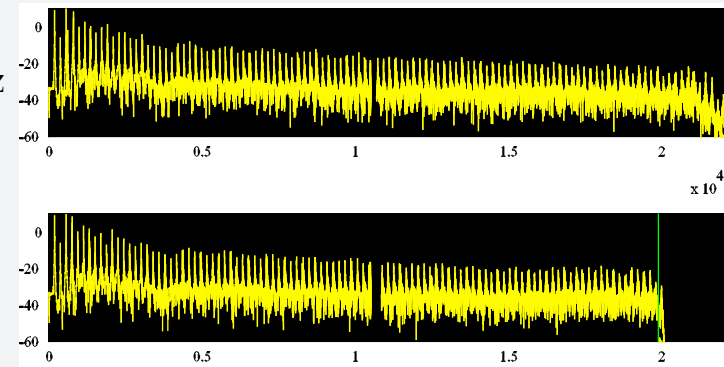
$$f_c = \frac{f_s}{\pi} \sin^{-1}(\lambda) \leq \frac{f_s}{2}$$

- Worst case near  $f_s = 44.1$  kHz, typical wind instrument dimensions:

$$f_c \approx 20 \text{ kHz}$$

$$f_s \approx 44037 \text{ Hz}$$

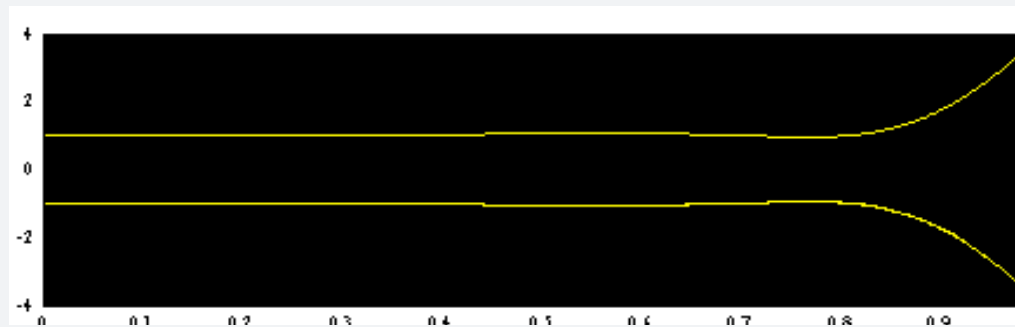
$$f_s \approx 44036 \text{ Hz}$$



# Accuracy—Modal Frequencies

- Numerical dispersion---normally a problem for FD schemes!
- This is a 2<sup>nd</sup> order scheme---might expect severe mode detunings...
- Not so...

E.g., for a lossless  
clarinet bore...



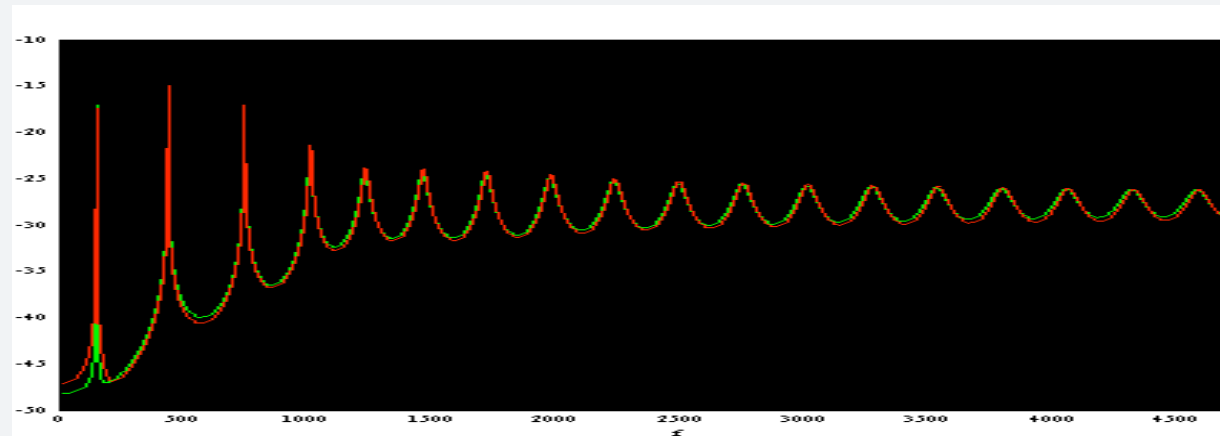
...calculated modal  
frequencies are nearly  
exact, over the entire  
spectrum

Mode #	Freq. (FD, Hz)	Freq. (exact, Hz)	cent diff.
1	141.89	141.96	0.86
2	413.79	413.95	0.65
3	705.55	705.55	0.00
12	3144.04	3142.63	- 0.77



# Accuracy—Transfer Impedance

- Even under more realistic conditions (i.e., with radiation loss), behaviour is extremely good:
- Transfer impedance (mouth  $\rightarrow$  radiating end):



**Red:** exact (calculated at 400 kHz)

**Green:** calculated at 44.1 kHz

- Upshot: FD approximation converges very rapidly...
- ...“perceptually” exact, even at audio sample rate.
- No compelling reason to look for better schemes...

# Explicit Updating

- Discretization of oscillator:

$$\ddot{y} + g\dot{y} + \omega_0^2 y + \omega_1^{1+\alpha} (-y - H)^+{}^\alpha = -a_1 p_\Delta$$

Parameterized implicit discretization

Implicit discretization

$$y^{n+1} - 2y^n + y^{n-1} + \frac{gk}{2}(y^{n+1} - y^{n-1}) + \boxed{\omega_0^2 k^2 \theta y^n + \frac{\omega_0^2 k^2 (1-\theta)}{2}(y^{n+1} + y^{n-1})} - \boxed{\frac{\omega_1^{\alpha+1} k^2}{2}(y^{n+1} + y^{n-1})(-y^n - H)^+{}^{\alpha-1}} = -a_1 p_\Delta^n$$

Exact integrator possible for linear part of oscillator...

Explicit update...

Mouthpiece pressure drop

$$p_\Delta^n = p_m^n - p_{in}^n$$

Flow nonlinearity

$$u_m^n = a_2 [y^n + H]^+ \sqrt{|p_\Delta^n|} \operatorname{sgn}(p_\Delta^n)$$

Flow conservation

$$u_{in}^n = u_m^n - u_r^n$$

Flow induced by reed

$$u_r^n = \frac{a_3}{2k} (y^{n+1} - y^{n-1})$$

Bore coupling

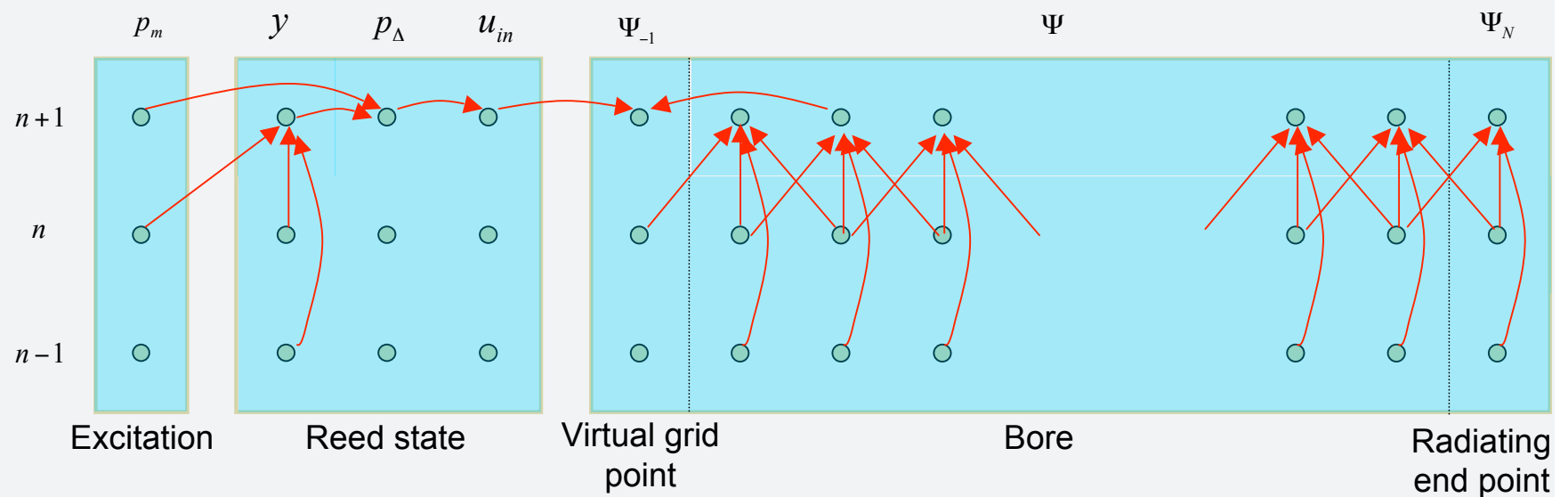
$$p_{in}^n = \frac{\rho}{2k} (\Psi_0^{n+1} - \Psi_0^{n-1})$$

$$u_{in}^n = -\frac{S_0}{2h} (\Psi_1^n - \Psi_{-1}^n)$$

- Implicit discretization → excellent stability properties
- Unknowns always appear linearly...

# Explicit Updating

- Can find a flow path in order to update all the state variables (sequentially)



- Similar to setting of “reflection-free port resistances” in linear WDF networks...
- ...but more general.

# Note on Stability

- The scheme for the bore + bell termination, in isolation, is guaranteed stable.
- Situation more complicated when reed mechanism is connected.
- Consider system under transient conditions (input  $p_m = 0$ ):

$$H(t) = H_{bore}(t) + H_{bell}(t) + H_{reed}(t) \leq H(0)$$

Total  
energy

Stored  
energy  
in bore

Stored  
energy  
at bell

Stored  
energy  
of reed

Initial  
stored  
energy

True...

- System is dissipative → state bounded for any initial conditions.
- Under forced conditions, would like:

$$H(t) \leq H(0) + \int_0^t f(p_m) dt'$$

Total  
energy

Initial  
stored  
energy

Energy  
supplied  
externally

Unfortunately  
this is false...

- Upshot: impossible to bound solutions of model system under forced conditions
- Best one can do: ensure energy balance is respected in FD scheme...

# Computational Cost

---

For a given sample rate  $f_s$ , bore length  $L$ , and wave speed  $c$ , the computational requirements are:

- $2Lf_s/c$  units memory
- $4Lf_s^2/c \rightarrow 6Lf_s^2/c$  flops/sec.

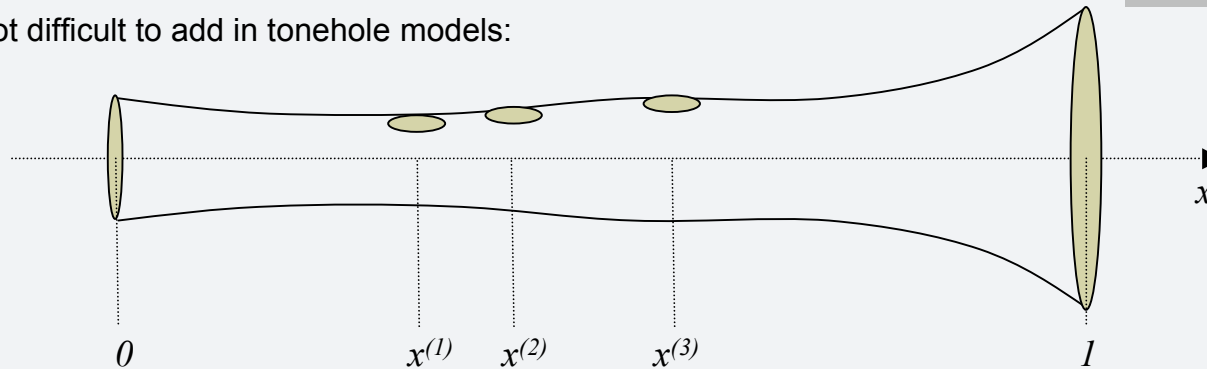
...independent of bore profile. Reed/tonehole/bell calculations are  $O(1)$  extra ops/memory per time step

Example: clarinet  $\rightarrow$  15 Mflops/sec., at  $f_s = 44.1$  kHz

Not a lot by today's standards...far faster than real time.

# Toneholes

- Not difficult to add in tonehole models:



- Can add terms pointwise to Webster's equation:

$$\Psi_{tt} = c^2 (S \Psi_x)_x + \sum_{q=1}^M \delta(x - x^{(q)}) m^{(q)}$$

$$m^{(q)} = A \Psi_{tt}(x^{(q)}, t) + B \Psi(x^{(q)}, t) \quad q = 1, \dots, M$$

State of tonehole  $q$

Physical parameters  
defining tonehole  $q$

- In FD implementation, can be added anywhere along bore (Lagrange interpolation):

# GUI: Matlab

**windfdgui**

**Global Simulation Parameters**

Sample Rate: 44100  
Duration (s): 3

**Tutorial Examples**

1 2 3 4 5  
6 7 8 9 10 11

**Reed parameters**

Surf. Area (m<sup>2</sup>): 0.000146  
Mass (kg): 3.372e-006  
Freq. (rad./s): 23250  
Damping (1/s): 3000  
Eq. disp. (m): 0.0004  
Chan. width (m): 0.01  
☒ reed beating on  
Beating param.: 300

**Tube Parameters**

wave speed (m/s): 340  
length (m): 0.664  
Bore profile type: cylinder with bell  
right/left area: 16.6  
bell fraction: 0.328

**Pressure Excitation**

ref. pres. (Pa): 2600  
rise time (s): 0.01  
ramp factor: 2  
decay time (s): 0.1

**Tremolo**

freq.(Hz): 3  
depth: 0.02

**Output**

Compute and Play Sound  
Play again  
Save as soundfile  
test

**Timer**

☐ Timer on

S. Bilbao, University of Edinburgh

**Toneholes**

☒ Toneholes on

	Pos (0-1)	rad. (m)	depth (m)
<input checked="" type="checkbox"/> Hole 1	0.6	0.005	0.003
<input checked="" type="checkbox"/> Hole 2	0.7	0.005	0.003
<input checked="" type="checkbox"/> Hole 3	0.8	0.005	0.003





**Events**

☒ Events on

	Time (s)	Hole #	final sta	duration
<input checked="" type="checkbox"/> Event 1	1	1	0	0.01
<input checked="" type="checkbox"/> Event 2	1.5	2	0	0.01
<input checked="" type="checkbox"/> Event 3	2	3	0	0.01

# Sound Examples

---

- Clarinet: 
- Saxophone: 
- Multiphonics/squeaks:  



# Conclusion

---

- Disadvantages:

- Costs more to compute than a typical waveguide model (but still not much!)

- Advantages:

- Bore modeling becomes trivial...
- More general extensions possible (NL wave propagation)
- Far more design freedom than, e.g., WG/WD methods